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# Structure and kinematics of the clouds surrounding the Galactic mini-starburst W43 MM1

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## ABSTRACT

Massive stars have a major influence on their environment, yet their formation is difficult to study as they form quickly in highly obscured regions and are rare, hence more distant than lower mass stars. Westerhout 43 (W43) is a highly luminous galactic massive star-forming region at a distance of 5.5 kpc and the MM1 part hosts a particularly massive dense core ( $1000 M_{\odot}$  within 0.05 pc). We present new *Herschel* HIFI maps of the W43 MM1 region covering the main low-energy water lines at 557, 987, and 1113 GHz; their  $\text{H}_2^{18}\text{O}$  counterparts; and other lines such as  $^{13}\text{CO}$  (10–9) and  $\text{C}^{18}\text{O}$  (9–8), which trace warm gas. These water lines are, with the exception of line wings, observed in absorption. *Herschel* SPIRE and JCMT  $450 \mu\text{m}$  data have been used to make a model of the continuum emission at the HIFI wavelengths. Analysis of the maps, and in particular the optical depth maps of each line and feature, shows that a velocity gradient, possibly due to rotation, is present in both the envelope ( $r \gtrsim 0.5$  pc) and the protostellar core ( $r \lesssim 0.2$  pc). Velocities increase in both components from SW to NE, following the general source orientation. While the  $\text{H}_2\text{O}$  lines trace essentially the cool envelope, we show that the envelope cannot account for the  $\text{H}_2^{18}\text{O}$  absorption, which traces motions close to the protostar. The core has rapid infall,  $2.9 \text{ km s}^{-1}$ , as manifested by the  $\text{H}_2^{18}\text{O}$  absorption features which are systematically redshifted with respect to the  $^{13}\text{CO}$  (10–9) emission line which also traces the inner material with the same angular resolution. Some  $\text{H}_2^{18}\text{O}$  absorption is detected outside the central core and thus outside the regions expected (from a spherical model) to be above 100 K; we attribute this to warm gas associated with the other massive dense cores in W43 MM1. Using the maps to identify absorption from cool gas on large scales, we subtract this component to model spectra for the inner envelope. Modeling the new, presumably corrected, spectra results in a lower water abundance, decreased from  $8 \times 10^{-8}$  to  $8 \times 10^{-9}$ , with no change in infall rate.

**Key words.** ISM: molecules – ISM: abundances – stars: formation – stars: protostars – stars: early-type – line: profiles

## 1. Introduction

Massive star formation is still poorly understood and while various means have been proposed to overcome the radiation pressure problem, observational ambiguities remain (see reviews by Zinnecker & Yorke 2007; Beuther et al. 2007). Massive dense cores are generally fairly distant and their centers always extremely embedded, making them difficult to observe with the necessary detail. In addition, the protostellar phase appears to be very short (a few  $10^5$  years, Tan et al. 2014; Beuther et al. 2007), further reducing the number of objects. A main question is how to “feed” a massive protostar, particularly in the earliest phases when it is weak in the infrared.

Understanding high-mass star formation is one of the goals of the Water In Star-forming regions with *Herschel*<sup>1</sup> program (WISH, Van Dishoeck et al. 2011). Water is a potentially very abundant molecule, but the vast majority of water is found as ice, which cannot be used to study kinematics. Only when the temperature reaches 100 K is the water ice on dust grains vaporized, provoking a radical change in water vapor abundance from some  $10^{-8}$  to  $1\text{--}10 \times 10^{-5}$  (Van der Tak et al. 2010;

Herpin et al. 2012). This jump in abundance makes it possible in principle to use the water molecule to focus on the central regions. High-energy transitions of other molecules also come from the dense and warm central regions. However, low-energy water transitions can be observed in both emission and absorption coming from the central regions, allowing potentially infalling cooler material to be detected in absorption, unlike in high-energy transitions (Van der Tak et al. 2013).

This idealized picture does not hold for the main water isotope as even the low-abundance gas is highly optically thick. However, for the  $\text{H}_2^{18}\text{O}$  and  $\text{H}_2^{17}\text{O}$  isotopomers with typical abundances of respectively 450 and 2000 times lower than the main isotope, the above picture should be more appropriate. The goal of this paper is to map W43 MM1 surroundings with a large set of lines to develop a more precise image of the amount of turbulence, infall rate, and rotation in the central core, but also in the cooler surrounding material. To this end, we try to identify the lines that best allow us to study the inner and outer parts of the protostellar envelope.

W43 MM1 is a high-mass proto-stellar object situated at about 5.5 kpc from the sun (Zhang et al. 2014) at the near end of the Galactic bar. W43 is sometimes considered a “mini-starburst” region because of the currently high level of star formation occurring there (Motte et al. 2003), probably due to a

<sup>1</sup> *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

**Table 1.** *Herschel*/HIFI mapped water line transitions towards W43 MM1.

Setting map (band)	Transition	Frequency [GHz]	$E_u$ [K]	Beam [ $''$ ]	Map size [ $'$ ]	$\eta_{\text{mb}}$	$T_{\text{sys}}$ [K]	rms [mK]
1 (1a)	o-H <sub>2</sub> <sup>18</sup> O 1 <sub>10</sub> –1 <sub>01</sub>	547.6764	60.5	37.8	0.92 × 1.74	0.62	80	27
1 (1a)	o-H <sub>2</sub> O 1 <sub>10</sub> –1 <sub>01</sub>	556.9361	61.0	37.1				
2 (4a)	p-H <sub>2</sub> O 2 <sub>02</sub> –1 <sub>11</sub>	987.9268	100.8	21.3	1.27 × 1.60	0.63	340	79
2 (4a)	C <sup>18</sup> O 9–8	987.5604	237.0	21.3				
3 (4b)	p-H <sub>2</sub> <sup>18</sup> O 1 <sub>11</sub> –0 <sub>00</sub>	1101.6982	52.9	19.9	1.35 × 1.6	0.63	390	20
3 (4b)	p-H <sub>2</sub> O 1 <sub>11</sub> –0 <sub>00</sub>	1113.3430	53.4	19.7				
3 (4b)	<sup>13</sup> CO 10–9	1101.3497	290.8	19.9				

**Notes.** Frequencies are from Pearson et al. (1991). The rms is the noise at  $\delta\nu = 1.1$  MHz.

cloud collision (Nguyen-Lu’o’ng et al. 2013). *Herschel* observations, including water lines, towards the center of W43 MM1 were analyzed by Herpin et al. (2012), who present a spherically symmetric model for the structure of W43 MM1 showing that both accretion and radially increasing turbulent velocities are required. They estimate the water abundances to be  $8 \times 10^{-8}$  where  $T_{\text{mb}} < 100$  K and  $1.4 \times 10^{-4}$  in the warm regions, higher than in most cores (Chavarría et al. 2010). The water abundance in cold outer envelopes is usually lower, on the order of  $10^{-9}$ – $10^{-10}$  (Van Dishoeck et al. 2013).

The general source structure is shown in Fig. 1 of Louvet et al. (2014): a double central core (N1a and N1b) surrounded by an elongated structure with a column density of  $\sim 10^{23}$  cm<sup>−2</sup>. Louvet et al. (2014) found a number of massive dense cores (MDC) in W43 MM1, mostly aligned to the NE and SW of the main core with a typical separation of about 5 to 10 $''$ . The maps presented in this work cover the region with the cluster of MDCs identified by Louvet et al. (2014, Fig. 1). The main MDC (N1) dominates the cluster, with an estimated mass of 2100  $M_{\odot}$ , with the other MDC masses decreasing to typically 100  $M_{\odot}$  away from the center of the cluster. All of the MDCs identified by Louvet et al. (2014) in W43 appear over-dense compared to other dense cores, such that they are likely to contain protostellar objects. Thus, the central spherically symmetric model may not accurately represent the source structure given the presence of heating sources – the MDCs – distributed over roughly an arcminute, particularly to the southwest of the center of W43 MM1.

As part of the WISH project, W43 MM1 was mapped with the HIFI heterodyne spectrometer aboard the *Herschel* Space Observatory (Pilbratt et al. 2010; Roelfsema et al. 2012) at 557, 987, and 1113 GHz. These frequencies correspond to the o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub>, p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub>, and p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> water lines, but also allow detection of the o-H<sub>2</sub><sup>18</sup>O 1<sub>10</sub>–1<sub>01</sub>, C<sup>18</sup>O 9–8, p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub>, and <sup>13</sup>CO 10–9 lines. The water lines (all isotopes) show a mixture of emission and absorption, while the CO isotopic lines are only seen in emission. Coupled with the spatial variations, this information is used to develop a more comprehensive model of the core and envelope structure and gas motions such as infall, turbulence, and/or rotation.

## 2. Observations

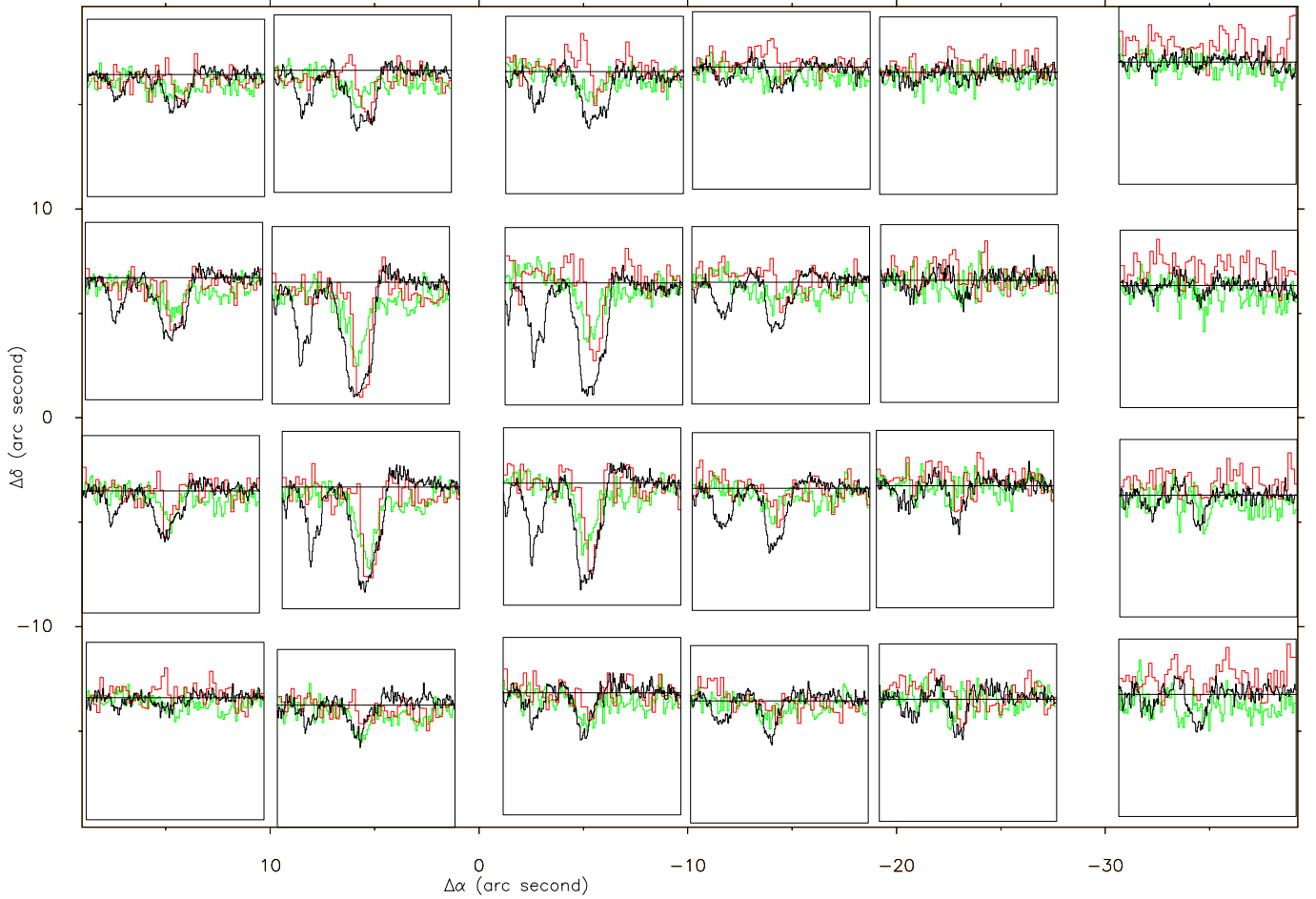
W43 MM1 was mapped at 557, 987, and 1113 GHz using the HIFI instrument (de Graauw et al. 2010), corresponding to the o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub>, p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub>, and p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> lines. The data were taken on October 27, 2010 (obsid 1342219193), March 11, 2011 (obsid 1342215899), and April 21, 2011 (obsid

1342207373). The *H* and *V* polarizations were observed simultaneously using both the acousto-optical Wide-Band Spectrometer (WBS) with 1.1 MHz resolution and the digital auto-correlator High-Resolution Spectrometer (HRS) at higher spectral resolution. We used the *On-The-Fly* mapping mode with Nyquist sampling and a reference position (RA<sub>OFF</sub> = 18h48min27s, Dec<sub>OFF</sub> = −1°44′28 $''$ ). The o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> line was mapped (map size 0.92′ × 1.74′) simultaneously with the o-H<sub>2</sub><sup>18</sup>O 1<sub>10</sub>–1<sub>01</sub> line in the lower sideband while the p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> line was mapped (1.27′ × 1.60′) with the C<sup>18</sup>O 9–8 in the same upper sideband. The third map (1.35′ × 1.68′) combined the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub>, p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub>, and <sup>13</sup>CO 10–9 lines. The (0, 0) position of all the maps is  $\alpha_{\text{J2000}} = 18^{\text{h}}47^{\text{m}}47.0^{\text{s}}$ ,  $\delta_{\text{J2000}} = -1^{\circ}54'28''$ . The observations are part of the WISH GT-KP (Van Dishoeck et al. 2011).

The three maps (see Table 1) were obtained using an *On-The-Fly* mapping procedure with spectra drifting along the declination axis on columns separated by a step chosen from the beam size of the main water line frequency. The pointing differs slightly between the *H* and the *V* receivers. For p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub>, the offset is of about 3 $''$  in right ascension, whereas the *On-The-Fly* right ascension step is 9 $''$ . As a consequence we have two maps shifted in right ascension by one-third of the right ascension step. For spectra with a low continuum level, the signal-to-noise ratio of the individual spectra is too low to have a good detection of the absorption features. In order to improve this signal-to-noise ratio, we then decided to produce an almost regular grid of averaged spectra by dividing the mapped area into cells on a regular grid. We averaged all spectra within a cell, both in intensity and position.

The off position has been inspected for each targeted frequency and does not show any emission. The frequencies, energy of the upper levels, system temperatures, integration times, and rms noise level at a given spectral resolution for each of the lines are provided in Table 1. Calibration of the raw data into the  $T_{\text{A}}$  scale was performed by the in-orbit system (Roelfsema et al. 2012); conversion to  $T_{\text{mb}}$  was done using the latest beam efficiency estimate from October 2014<sup>2</sup> given in Table 1 and a forward efficiency of 0.96. HIFI receivers are double sideband with a sideband ratio close to unity (Roelfsema et al. 2012). The flux scale accuracy is estimated to be 10% for band 1 and 15% for band 4<sup>2</sup>. The frequency calibration accuracy is 20 kHz and 100 kHz respectively for HRS and WBS observations. Data calibration was performed in the *Herschel* Interactive Processing Environment (HIPE, Ott 2010) version 13. Further analysis was

<sup>2</sup> <http://herschel.esac.esa.int/twiki/bin/view/Public/HifiCalibrationWeb?template=viewprint>



**Fig. 1.** Map of observed spectra towards W43 MM1 for lines p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> (black,  $T_{\text{mb}} = [-2.2, 1.00]$  K), p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> (red,  $T_{\text{mb}} = [-0.55, 0.25]$  K), and <sup>13</sup>CO (10–9) (green,  $T_{\text{mb}} = [1.0, -0.50]$  K, inverted  $T$  scale in order to ease comparisons). Velocity range:  $V_{\text{lsr}} = [70, 130]$  km s<sup>–1</sup>. Plotted spectra are an average of *On-The-Fly* individual spectra around each box central position. The continuum level is subtracted.

done within the CLASS<sup>3</sup> package. These lines are not expected to be polarized; thus, after inspection, data from the two polarizations were averaged together.

### 3. Mapping result

The observations discussed here map, in three different frequency ranges, a region of about 80'' by 80'' centered on the W43 MM1 core observed, with a very good signal-to-noise ratio, in Herpin et al. (2012). With a much higher rms per spectrum, our maps still exhibit the same velocity components, in varying proportions over the map: a wide ( $\Delta V \simeq 40$  km s<sup>–1</sup>) emission, a narrower ( $\Delta V \simeq 10$  km s<sup>–1</sup>) emission, and a rather deep absorption of a continuum background component. Figures 1–3 show the spectra around MM1 for a velocity range bracketing the main water line and also, in Figs. 1 and 3, the highest velocity diffuse line of sight cold cloud (see Herpin et al. 2012). Strongly dominated by the MM1 core, these maps because of an extended continuum background still reveal structure far from MM1, making the continuum component a key feature to trace the water lines outside the MM1 core.

The analysis of Herpin et al. (2012) needed a cold water cloud to correctly describe the various water line profiles. Unlike the core, this foreground material might be resolved in these

maps which show that p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> is mainly detected around MM1, while the lines of p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> and o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> are detected at all positions having a positive continuum level. The continuum background (Fig. 4) which peaks at 2.5 K for p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> at the W43 MM1 core position is almost completely absorbed for channels close to the source velocity. However, both blue and red wings of the main absorption showing different behaviors around the core help analyze the cold foreground material close to MM1.

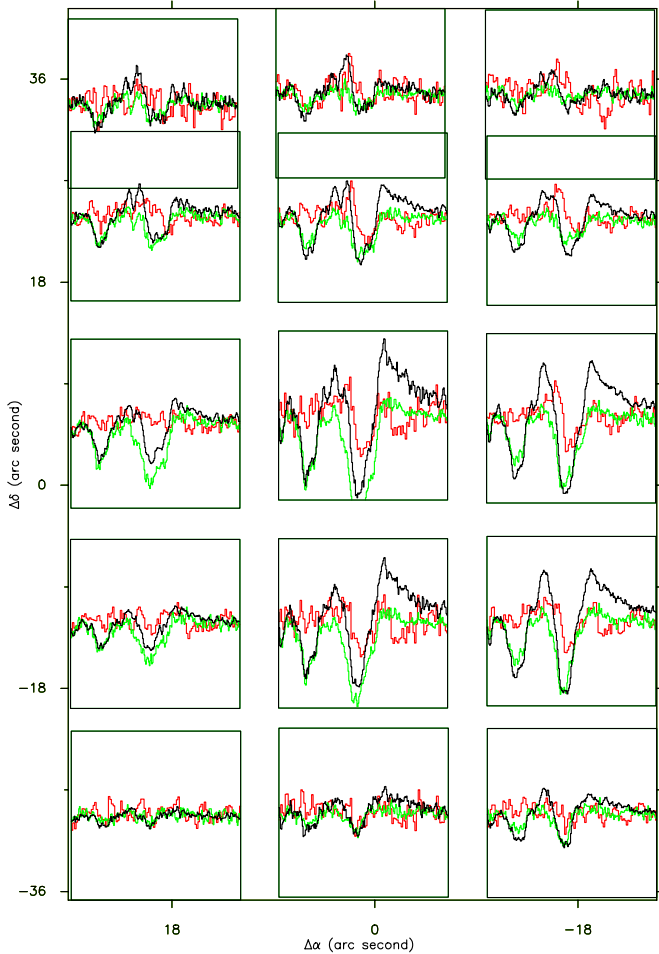
Figure 1 shows the observed spectra towards W43 MM1 for lines p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> (black), p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub>, and <sup>13</sup>CO (10–9) (inverted to help comparison). These water spectra exhibit a strong absorption against an extended continuum. The HIFI beam does not resolve the MM1 core material above 100 K, but it is small enough to map the continuum region and hence the very close surroundings of MM1.

Figure 2 shows the observed spectra towards W43 MM1 for lines o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub>, p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub>, and o-H<sub>2</sub><sup>18</sup>O 1<sub>10</sub>–1<sub>01</sub>. The water spectra are seen in absorption and are thus mainly seen toward the MM1 core.

Figure 3 shows the observed spectra towards W43 MM1 for lines p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> and C<sup>18</sup>O (9–8) (restricted to spectra with a sufficient signal-to-noise ratio). The angular resolution is similar to the map size and help very little to map the MM1 surrounding material.

<sup>3</sup> <http://www.iram.fr/IRAMFR/GILDAS/>





**Fig. 2.** Map of observed spectra towards W43 MM1 for lines o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> (black,  $T_{\text{mb}} = [-0.5, 0.5]$  K), p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> (green,  $T_{\text{mb}} = [1.0, 1.0]$  K), and o-H<sub>2</sub><sup>18</sup>O 1<sub>10</sub>–1<sub>01</sub> (red,  $T_{\text{mb}} = [-0.15, 0.15]$  K). Velocity range:  $V_{\text{lsr}} = [70, 130]$  km s<sup>−1</sup>. The plotted spectra are an average of *On-The-Fly* individual spectra around each box central position. The continuum level is subtracted. Axis units are offset (″) from MM1.

### 3.1. Continuum background

Though our spectra do contain continuum, its precise level is somewhat uncertain. We mapped the water lines in W43 MM1 using a fast *On-The-Fly* mapping HIFI-AOT which saved observing time, but at a price. The continuum level between two consecutive spectra along the drifts exhibits variations much above those expected from the spectra rms. These variations are especially large at the edges of the region where the continuum level is weaker. One even detects lines in absorption, though the continuum level appears negative in the spectrum. Mapping the continuum derived from our maps produces various small artifacts due to these random level variations. The overall structure of the region is correctly reproduced, but absolute values of the continuum are unreliable at the smaller scales. As we want to use the continuum level to evaluate the opacity of the absorbing water material, it will produce large errors in parts of the map of low continuum intensity. To overcome this problem we decided to estimate the true value of the continuum level by using data obtained with SPIRE at 250 μm (Nguyen-Lu’o’ng et al. 2013). Our best map for comparison is the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> map whose wavelength is 269 μm, quite close to the SPIRE data. Figure 4 shows an image of the continuum level observed by SPIRE at 250 μm whose angular resolution was smoothed to the p-H<sub>2</sub>O

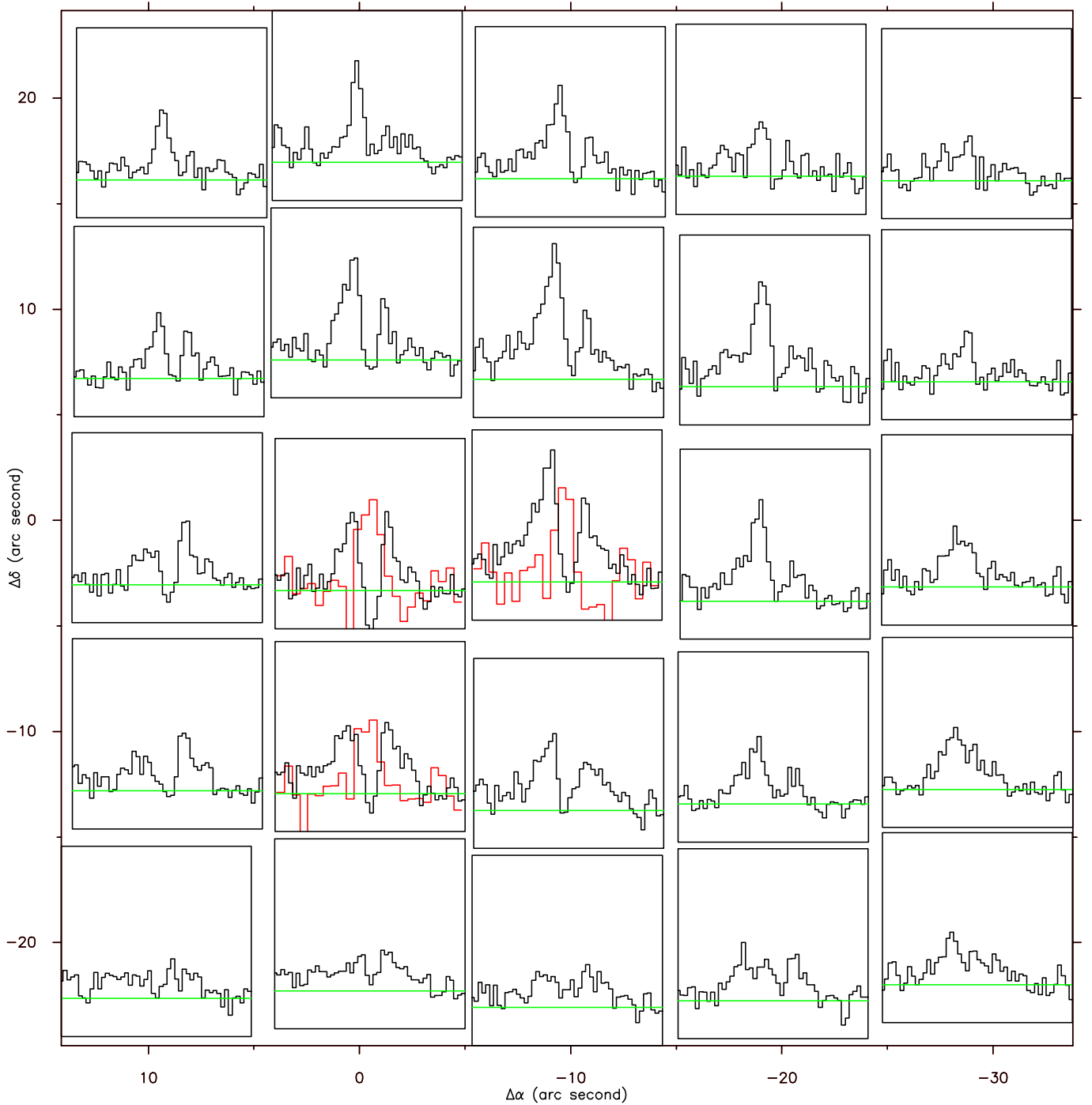
1<sub>11</sub>–0<sub>00</sub> resolution. The continuum derived from the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> spectra observed at 1113 GHz is overplotted as contour levels. The SPIRE map of the continuum has a strong and unresolved peak located at the MM1 position, but at lower levels it follows an underlying structure aligned northeast to southwest and anchored to the MM1 position. It decreases rapidly towards the edges of the map along the southwest map axis and more slowly towards the southeast part of the map. The smallest beam size ( $\approx 20''$ , p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub>) does not resolve the  $10''$  central MM1 source, but does resolve the extended continuum structure. Thus, having observed absorption lines more than a beam away from MM1 we can get an insight into the material surrounding the central core. Despite level uncertainties in the HIFI map, a comparison of both maps shows a small offset ( $-4''$ ,  $+2''$ ) between the two peak locations. The reason for this discrepancy is unknown. A map of the region obtained by SCUBA/JCMT at 450 μm (Chavarria, priv. comm.) is also available. The SPIRE and SCUBA maps have a very similar geometry apart from a less than  $3''$  north–south shift between the peak positions. Part of the data analysis which follows uses, for p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub>, the computation of the absorbed fraction  $\frac{T - T_{\text{cont}}}{T_{\text{cont}}}$ . The error on this quantity is highly dependent on the  $T_{\text{cont}}$  value. The low reliability on the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> continuum level makes it too difficult to use. Thus, we decided to modify the observed continuum level in all spectra using a model derived from the SPIRE map.

### 3.2. Model of the continuum background

We build a new set of p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> spectra by setting the continuum level in each spectrum to a new value derived using the following procedure. We build the continuum model of the region at 1113 GHz from the SPIRE map smoothed at the HIFI p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> spatial resolution to obtain the map geometry. We then derive the intensity of the peak from the Herpin et al. (2012) observation. For the latter, one could note that the main absorption in this spectrum goes below zero. This observation was done in a beam switch mode. The SPIRE map showed that the OFF positions were not fully free of continuum, especially the east one. Taking this into account – as well as the spectral index between the USB and LSB band and the slight offset between this spectrum and the position of the maximum in the map – we set the map maximum to 3.71 K. The second step of our continuum building process is to shift the HIFI map by ( $-4''$ ,  $+2''$ ) in order to correct the observed discrepancy between maps (see above). We think that this step is necessary, otherwise in the final map many of the spectra having the deepest absorption will fall below the zero level. Once done, we can easily find the new continuum level value for each spectrum at the corresponding offset in the model.

### 3.3. Spectral maps

All data exhibit various absorption features between 0 km s<sup>−1</sup> and 140 km s<sup>−1</sup>. At almost all map positions, even further than one p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> beam size, all water lines, as well as p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub>, exhibit an often wide and well-defined absorption close to the source velocity. The detection of weaker and narrower absorptions at positions distant from MM1 by much more than a beam size proves that the absorbing material extends over a region bigger than the unresolved MM1 core itself. Many narrower absorptions are also seen at lower velocities ranging from 85 km s<sup>−1</sup> down to 0 km s<sup>−1</sup>. These were interpreted by Herpin et al. (2012) as absorbing material on the line of sight



**Fig. 3.** Map of observed spectra towards W43 MM1 for lines p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> (black,  $T_{\text{mb}} = [-0.3, 1.2]$  K) and C<sup>18</sup>O (9–8) (red,  $T_{\text{mb}} = [-0.1, 0.4]$  K). We only plotted the three C<sup>18</sup>O (9–8) spectra with a good enough signal-to-noise ratio. Velocity range:  $V_{\text{lsr}} = [70, 130]$  km s<sup>−1</sup>. Plotted spectra are an average of *On-The-Fly* individual spectra around each box central position. The continuum level is subtracted. Axis units are offset (″) from MM1.

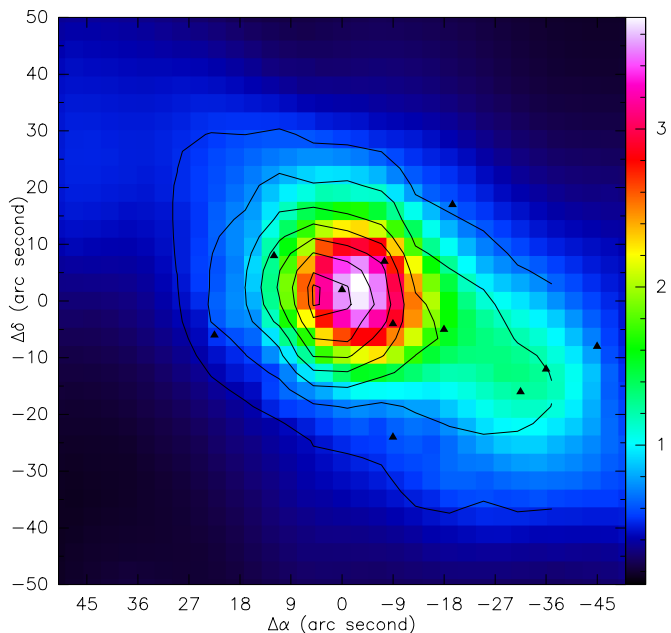
in front of the continuum background, the so-called diffuse line of sight cold clouds. This cold material is thought to be spatially unrelated to MM1 (see Van der Tak et al. 2013).

The three maps discussed in this paper are of unequal interest to investigate the MM1 region structure. The o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> spectra have a good rms but an angular resolution about four times larger than the source size. The p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> map contains a few bad spectra for one of the *On-The-Fly* passes, giving rise to a high signal-to-noise. The p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> line, detected in absorption against the extended continuum, with the best rms

among the three maps, is thus the most promising line to map the kinematics and physical structure at sizes comparable to the MM1 core. Opacity maps later derived from these data will help to separate line components over the map (see Figs. 5 and 6).

### 3.3.1. Main absorption component

At the MM1 position, the absorption velocities of the three water lines vary between 97.8 and 99.3 km s<sup>−1</sup>, whereas o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> are seen strongly absorbed at the source



**Fig. 4.** Maps of the W43 MM1 continuum level at 250  $\mu\text{m}$  (SPIRE data) and at 1113 GHz ( $\approx 269 \mu\text{m}$ ) before the offset correction. HIFI contour levels: 0.6, 1.0, 1.5, 2.0, 2.5, 3.0, 3.5, 3.65 K with original data positions but after continuum level correction (see text). The SPIRE data (in K) resolution is smoothed to the  $20''$  beam size of p-H<sub>2</sub>O  $1_{11}-0_{00}$ . Black triangles indicate the Louvet et al. (2014) sources.

velocity. The p-H<sub>2</sub>O  $2_{02}-1_{11}$  line has a higher energy level and its profile exhibits a clearly visible emission component (see central position of Figs. 1–3). As is expected, the deepest p-H<sub>2</sub>O  $1_{11}-0_{00}$  absorption in the map is indeed observed at the location of the strongest continuum component, i.e., the MM1 location.

### 3.3.2. p-H<sub>2</sub>O $1_{11}-0_{00}$ map

Figure 1 shows spectra for p-H<sub>2</sub>O  $1_{11}-0_{00}$ , p-H<sub>2</sub><sup>18</sup>O  $1_{11}-0_{00}$ , and <sup>13</sup>CO (10–9). The <sup>13</sup>CO (10–9) line is plotted inverted to ease the profile comparisons. The spectral resolution differs for each line. It is clear that the p-H<sub>2</sub><sup>18</sup>O  $1_{11}-0_{00}$  spectra (in red) are narrower than the p-H<sub>2</sub>O  $1_{11}-0_{00}$  spectra and located on the red part of the p-H<sub>2</sub>O  $1_{11}-0_{00}$  line, whereas the <sup>13</sup>CO (10–9) line (in green) appears more blueshifted. It can also be seen that the diffuse line of sight cold clouds are detected only for p-H<sub>2</sub>O  $1_{11}-0_{00}$  and o-H<sub>2</sub>O  $1_{10}-1_{01}$ , as already stated in Herpin et al. (2012).

### 3.3.3. o-H<sub>2</sub>O $1_{10}-1_{01}$ map

Figure 2 shows spectra toward W43 MM1 for o-H<sub>2</sub>O  $1_{10}-1_{01}$  and o-H<sub>2</sub><sup>18</sup>O  $1_{10}-1_{01}$ . The p-H<sub>2</sub>O  $1_{11}-0_{00}$  line spectrum is added for comparison of the maps. Despite a  $40''$  beam size comparable to the map size, there are various differences between the two lines over the maps.

The diffuse line of sight cold cloud absorptions are detected in o-H<sub>2</sub>O  $1_{10}-1_{01}$ , but not in the o-H<sub>2</sub><sup>18</sup>O  $1_{10}-1_{01}$  isotopic line, except for the  $(0'', +35'')$  position. The diffuse line of sight cold cloud signal evolution correlates strongly over the map with the corresponding signal in the p-H<sub>2</sub>O  $1_{11}-0_{00}$  line.

Both the o-H<sub>2</sub>O  $1_{10}-1_{01}$  line (black spectrum) and the o-H<sub>2</sub><sup>18</sup>O  $1_{10}-1_{01}$  (red spectrum) line are clearly detected over the whole map. The o-H<sub>2</sub>O  $1_{10}-1_{01}$  line appears as the merging of a deep absorption and of a wider emission component only seen

for the closest spectra to MM1. The absorbing component appears narrower for o-H<sub>2</sub>O  $1_{10}-1_{01}$  than for the p-H<sub>2</sub>O  $1_{11}-0_{00}$  line (green spectrum) at the central position and much more similar in the outer part of the map. The wide emission is brighter than for p-H<sub>2</sub>O  $1_{11}-0_{00}$ .

The o-H<sub>2</sub><sup>18</sup>O  $1_{10}-1_{01}$  main absorption appears narrower than o-H<sub>2</sub>O  $1_{10}-1_{01}$  and is also redshifted.

### 3.3.4. p-H<sub>2</sub>O $2_{02}-1_{11}$ map

Figure 3 shows spectra for p-H<sub>2</sub>O  $2_{02}-1_{11}$  and C<sup>18</sup>O (9–8). This *On-The-Fly* map was achieved in two passes. In the second one, a few groups of spectra suffered from poor baselines making the resulting map inhomogeneous if using the summed spectra, or the rms too low if using only the first pass. Figure 3 shows the sum of the two passes. The C<sup>18</sup>O (9–8) spectra are the most affected by the bad baseline, and so the central spectrum is the only one shown in Fig. 3. The C<sup>18</sup>O (9–8) line is seen in emission at the velocity of the absorbed p-H<sub>2</sub>O  $2_{02}-1_{11}$  component.

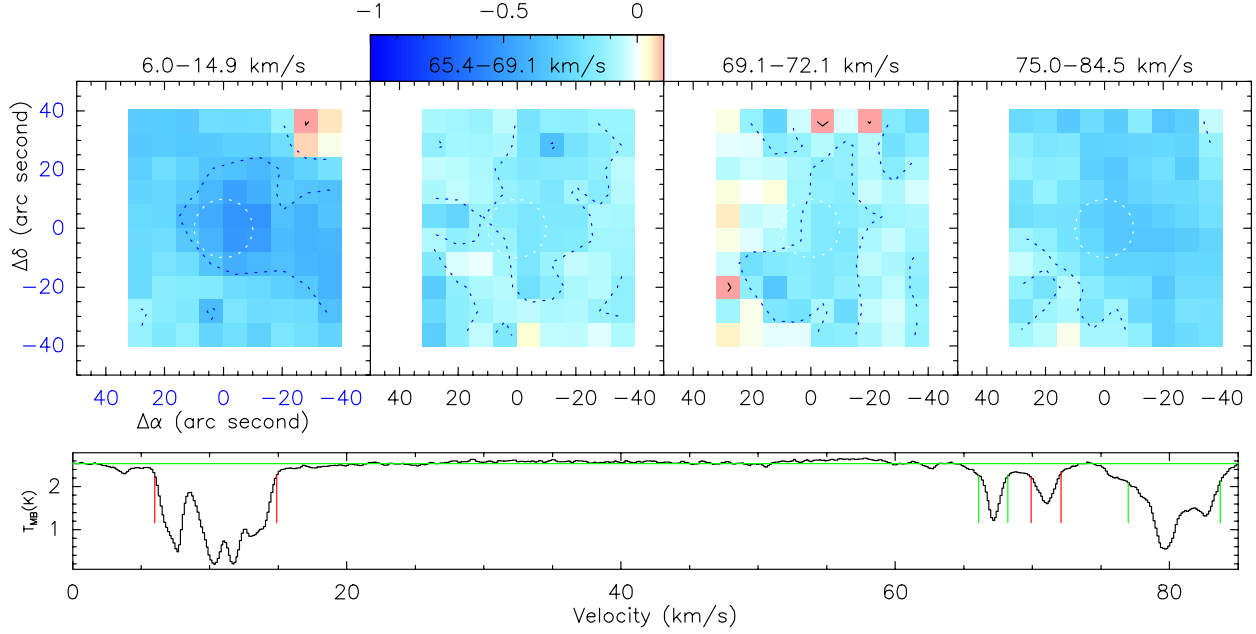
Though the high rms of these spectra makes it difficult to detect p-H<sub>2</sub>O  $2_{02}-1_{11}$  lines that are too weak, we do detect these lines off the MM1 center. The  $21.3''$  beam size is small enough to allow some mapping off the source core, in particular of the wings. It is clear that the line profile changes around the spectrum at the  $(0, 0)$  position. The blue wing part of the emission is seen up to the  $(0'', 20'')$  position, whereas the red part has almost disappeared. The red wing is more pronounced southwest of MM1. There is also a reliable signal in the southeast part of the map, see the  $[-30'', -20'' \text{ to } 0'']$  three spectra. We also note that the red wing is still detected to the west of MM1 whereas the blue wing became weaker.

### 3.4. Absorbed fraction maps

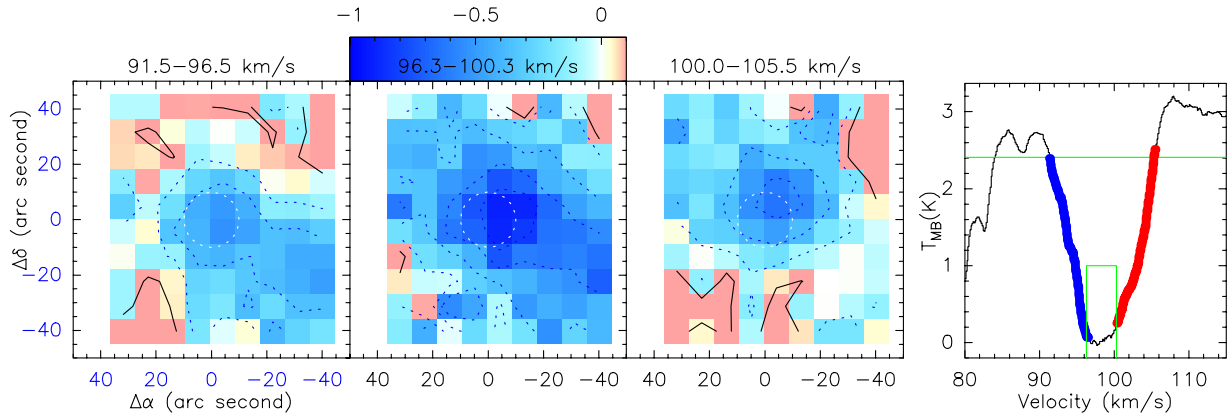
As explained above, variations in the intensities of the absorptions are dominated by the variations of the continuum background. Hence, mapping these intensities does not map water, but mostly reproduce the continuum geometrical structure over the map. However, it is possible to remove this continuum contribution by plotting the absorbed fraction  $\frac{T-T_{\text{cont}}}{T_{\text{cont}}}$  over the mapped region. Figures 5 and 6 show the integrated value of the absorbed fraction within a selected velocity range. Foreground clouds detected in absorption were identified in Chavarría et al. (2010) and Herpin et al. (2012) from their velocities. Here we use a different method based on the idea that foreground clouds are necessarily nearer and thus likely to be fairly homogeneous over the small scales (sub-pc) observed in the maps presented in this article. Figure 5 shows the behavior of absorbed fraction over the map for the four main diffuse line of sight cold clouds. The MM1 source is within the white dashed circle. We note that the absorbed fraction is almost constant at all positions following the continuum background. This favors the idea that the signal originates from some foreground material unconnected with MM1 itself.

Figure 6 plots the absorbed fraction for three parts of the main p-H<sub>2</sub>O  $1_{11}-0_{00}$  absorption, namely the blue wing of the line, its central part, and its red wing. A different behavior can be seen here for the diffuse line of sight cold clouds. Only the outer part of the MM1 core can be seen; the core material itself is shielded. In addition, the blue and red wings behave differently as they are located symmetrically around MM1.

Figure 7 compares absorbed fraction spectra computed for p-H<sub>2</sub>O  $1_{11}-0_{00}$  and p-H<sub>2</sub><sup>18</sup>O  $1_{11}-0_{00}$ . The noise in absorbed



**Fig. 5.** Maps of the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> average absorbed fraction  $\langle \frac{T - T_{\text{cont}}}{T_{\text{cont}}} \rangle$  in four velocity ranges bracketing the main four diffuse line of sight cold clouds. Velocity ranges are given above each box and are marked on the low signal-to-noise spectrum observed at (0, 0) by Herpin et al. (2012), the horizontal green line marks the continuum level (in K). The color scale is clipped in the range [−1, +0.1]. The white dashed circle stands for the beam at 1113 GHz at the (0, 0) position. These maps show an almost constant signal over all positions having a detectable continuum level. The MM1 core, located inside the beam circle is barely visible.



**Fig. 6.** From left to right: maps of the average absorbed fraction  $\langle \frac{T - T_{\text{cont}}}{T_{\text{cont}}} \rangle$  of the main p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> line blue wing, center, and red wing. Right: the (0, 0) spectrum as in Fig. 5 (in K), the blue and red wing ranges are color-coded. Other plot options are the same as in Fig. 5. These maps show that the blue wing material peaks southeast, whereas the red wing peaks north-northeast.

fraction spectra increases when  $T_{\text{cont}}$  decreases making the outer part of the map too noisy. To handle this difficulty we applied a variable spectral resolution to these spectra in order to keep this noise within acceptable values. The spectral resolution decreases with  $T_{\text{cont}}$ , mainly with the distance to MM1. After subtraction of the continuum contribution to the geometry of the map, these spectra present the spatial variation of the absorbing material. A similar absorbed fraction means a similar opacity. This figure confirms that the blue wing of p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> has almost no corresponding p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> counterpart unlike the red one.

#### 4. Velocity structure of the W43 core

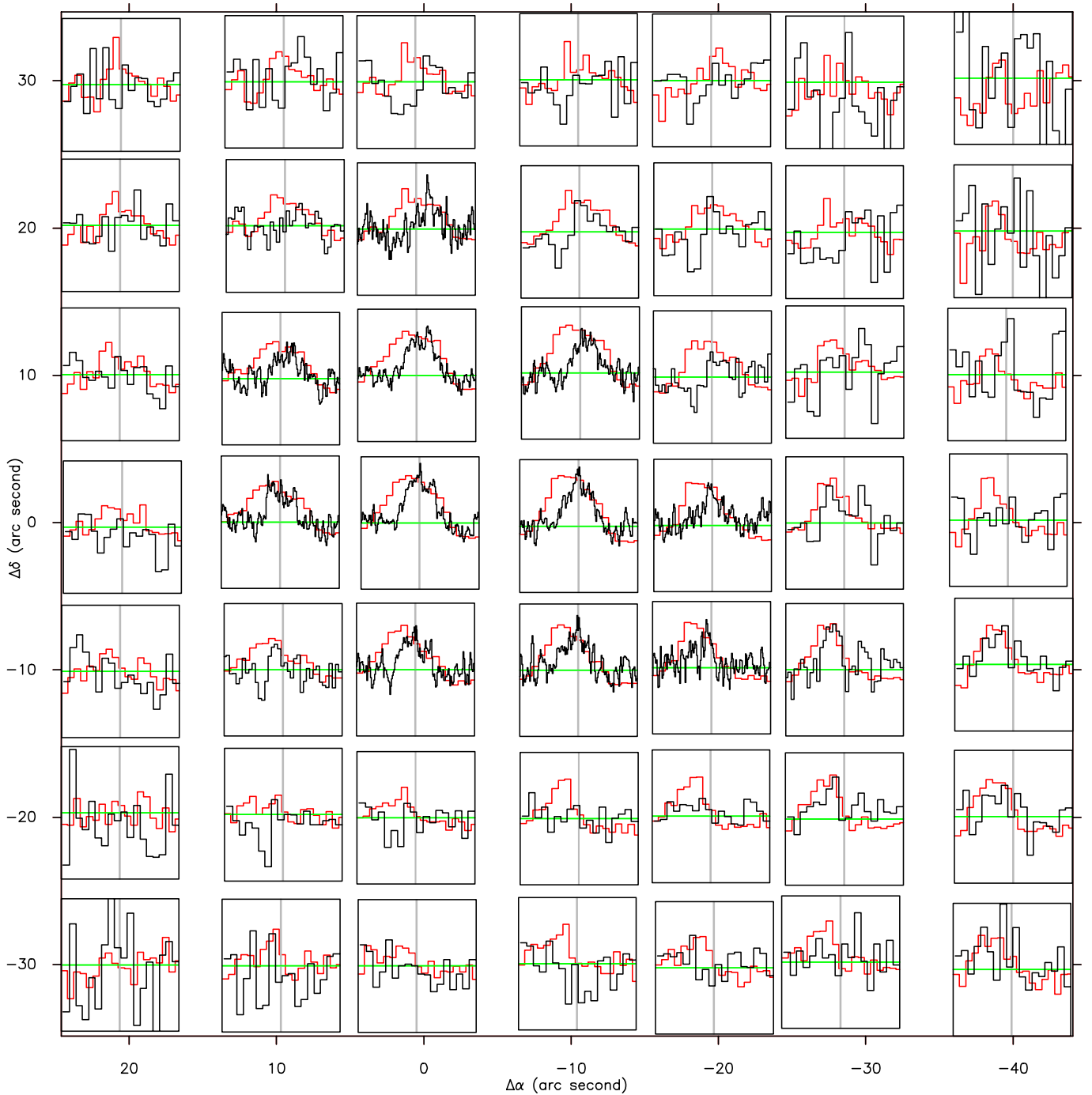
The comparison of the absorptions of the main H<sub>2</sub>O water lines (p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> and o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub>) and of the H<sub>2</sub><sup>18</sup>O lines (p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> and o-H<sub>2</sub><sup>18</sup>O 1<sub>10</sub>–1<sub>01</sub>) (see Figs. 1, 2, and 7) shows that, at all the detected positions, the blue side wing of

the main water line has no H<sub>2</sub><sup>18</sup>O counterpart. Hence, we think that this blue side of the main water absorption is due to a cold foreground material, possibly surrounding the MM1 core. An outflow absorption is partly taken into account in the models in Sect. 6.2, but we cannot rule out a low density outflow as the source of this absorption. Figure 6 also shows that the redshifted material is located at the immediate northeast off the MM1 core and that the blueshifted material is located southwest of MM1. The central and deepest part of the absorption covers both regions. Though the H<sub>2</sub><sup>18</sup>O absorption probably comes from the inner region of the source, it is also seen off-core towards the southwest MM1 following the extended continuum emission.

##### 4.1. Central velocity gradient

Due to its lower optical depths, p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> traces the inner material whereas the main isotope is optically thick in the

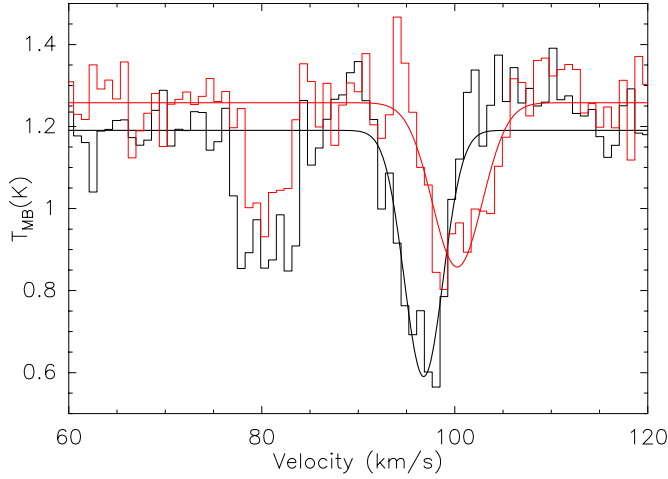




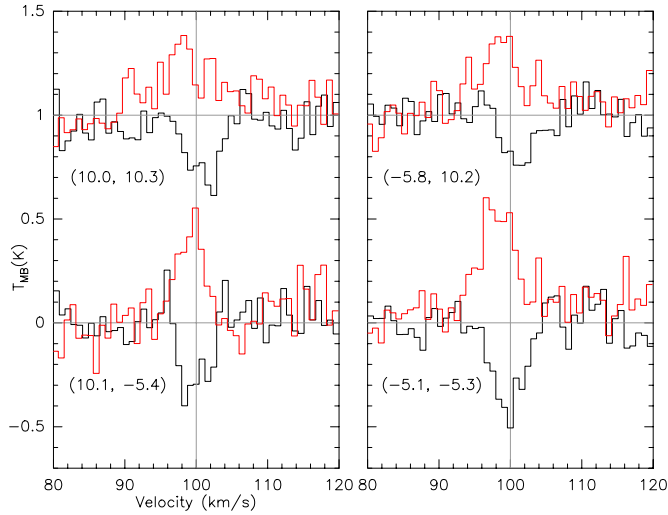
**Fig. 7.** Map toward W43 MM1 of  $\frac{T_{\text{cont}} - T}{T_{\text{cont}}}$  for p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> (red) and p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> (black) in the velocity range [90 km s<sup>−1</sup>, 110 km s<sup>−1</sup>]. The y-axis is from −0.1 to 1.1 for p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub>, and p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> is scaled by 4. The horizontal green lines define the zero level. The vertical gray lines mark 100 km s<sup>−1</sup>. As the signal-to-noise ratio increases for decreasing  $T_{\text{cont}}$ , the spectra velocity resolution varies over the map in order to keep it at a minimum value. Hence, with the greatest  $T_{\text{cont}}$ , central spectra having a greater number of channels appear plotted darker than those with a lower  $T_{\text{cont}}$ . Axis units are offset (″) from MM1.

surrounding cool gas, such that it does not trace the inner material. Figures 8 and 9 show respectively the main isotope and the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> ground state absorption lines at selected positions around MM1. Figure 8 compares the velocity profiles of p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> spectra between two positions, chosen symmetrically relative to MM1 and separated by more than a full beamwidth, to the SW and NE of MM1. These positions were chosen to be along the extended continuum object, far enough not to detect the emission from the central source but in a region

where the absorption is still clearly detected. It is immediately apparent that the W43 MM1 absorption shows a velocity gradient whereas the diffuse line of sight cold cloud at 80 km s<sup>−1</sup> does not. The diffuse line of sight cold clouds absorptions are rather identical in absorbed fraction where the main absorption shows a kind of complementary profiles. The blue wing is in emission for one position and in absorption for the other one. The opposite behavior is seen for the red wing. This inversion occurs, in velocity, symmetrically to the ≈99 km s<sup>−1</sup> source velocity.



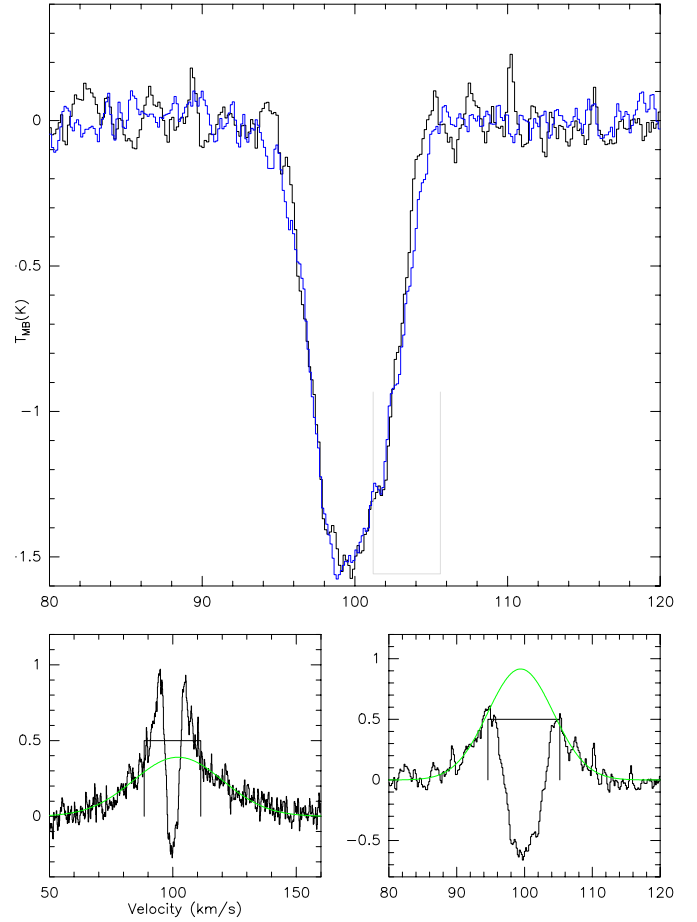
**Fig. 8.** Comparison of p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> spectra towards our map at offsets around (–18'', –20'') in black and around (+18'', +20'') in red.



**Fig. 9.** p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> averaged spectra around the (0, 0) position. The average center position is given in '' offsets below each spectrum. Top spectra were shifted by 1 K for plotting purposes. Averaged <sup>13</sup>CO (10–9) spectra (same positions) are overplotted in red. The continuum is removed for both lines, horizontal gray lines mark the zero levels, and the vertical gray lines mark 100 km s<sup>–1</sup>.

The p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> is not detectable as far out but Fig. 9 shows a similar velocity gradient in the same direction as for the envelope. It shows four p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> spectra at positions surrounding MM1 revealing a velocity gradient in the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> ranging from  $\approx 98$  km s<sup>–1</sup> in the southwest of MM1 to  $\approx 102$  km s<sup>–1</sup> in the northeast part of the MM1 core. The vertical line at the central velocity makes it clear that there is a SW-NE gradient but not SE-NW.

The <sup>13</sup>CO (10–9) emission line (Figs. 1, 9), whose energy is above 200 K, comes from the warm central regions. It has the same gradient but the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> absorption is consistently redshifted with respect to the <sup>13</sup>CO (10–9), showing that the infall onto the main source is visible throughout the inner region. As for the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> line, the <sup>13</sup>CO (10–9) line velocity shift of 0.7 km s<sup>–1</sup> is seen from southwest (offset –8'', –8'') to northeast (offset +8'', +8''), going from 97.9 to 98.6 km s<sup>–1</sup>. While the size of our beam and the signal-to-noise level of the observations does not allow us to measure the size of the region

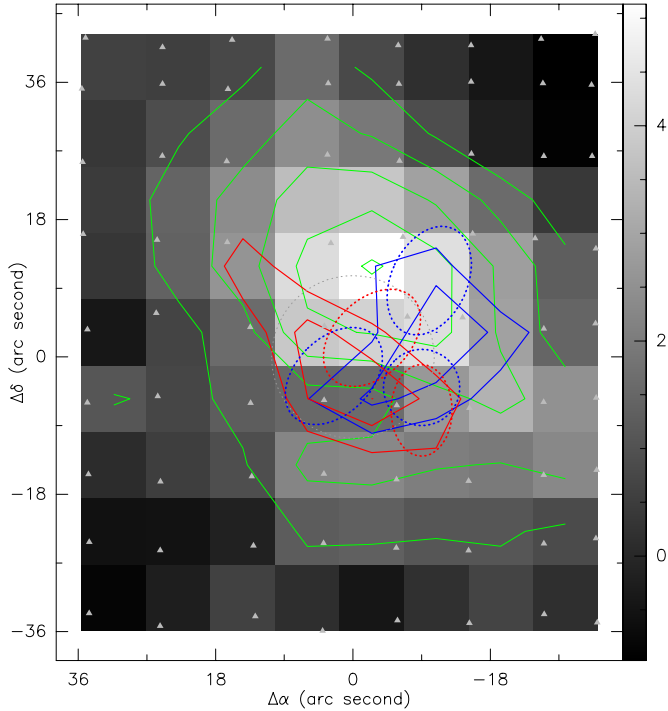


**Fig. 10.** Absorption of p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> versus p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> (continuum level subtracted) after removal of a wide component and then a narrow Gaussian component. The wide component is fitted in the *lower left box*. The removed wide spectrum is shown in the *lower right box* with a fit of the narrow component. The *upper box* shows the final result where the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> spectrum intensity is scaled to the absorption level of p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub>. This shows the high similarity of both profiles. Intensities are in K in all boxes.

from which <sup>13</sup>CO (10–9) emission is detected, the shift shows that it is not point-like even to our 20'' beam. The small velocity shift shows that the core is not rotationally supported, in agreement with the presence of the infall signature.

#### 4.2. p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> versus p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub>

Figure 10 compares the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> profile to the p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> absorption profile at the MM1 position. We first fit a Gaussian profile to model the wide component (see lower left spectrum in figure). Once done, we removed it from the spectrum. Then, we fit the emission part of the remaining spectrum by another Gaussian profile (see lower right spectrum in figure) that we again remove. The hypothesis here is that both the wide and the narrower emission parts of the line profile can be modeled by Gaussian profiles. The last spectrum fits then the absorbing component in the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> line. We compare its profile to the p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub> line by scaling both lines to the same intensities (see upper box in figure). Both profiles are almost identical in width. The little dip in the red wing is even the same in the two spectra. We conclude that both lines originate from the same material.



**Fig. 11.** Map of the blue and red wings of the wide  $p\text{-H}_2\text{O } 2_{02}-1_{11}$  component. Integrated line areas are in velocity ranges ( $\text{km s}^{-1}$ ): blue = [86, 92], red = [107, 117], green = [92, 107] (line center, also plotted as gray background). The gray dashed circle at the center gives the beam size. Dashed ellipses mark the position of the three CO (3–2) outflows (OF-1 and OF-2 red TBN) of Fig. 2 in Sridharan et al. (2014).

#### 4.3. The wide component high-velocity gas

A wide velocity component is observed in several main water lines (see Herpin et al. 2012). We detect it in our maps of  $o\text{-H}_2\text{O } 1_{10}-1_{01}$ ,  $p\text{-H}_2\text{O } 2_{02}-1_{11}$  and  $p\text{-H}_2\text{O } 1_{11}-0_{00}$ . The wide component in the  $o\text{-H}_2\text{O } 1_{10}-1_{01}$  map is detected at all observed positions but not used in our discussion because of the large beam size at 557 GHz. Because of a poor rms and of some baseline problems there are large uncertainties for the  $p\text{-H}_2\text{O } 2_{02}-1_{11}$  spectra. Still, it is possible to map the wide  $\approx 40 \text{ km s}^{-1}$  component wings of  $p\text{-H}_2\text{O } 2_{02}-1_{11}$  as seen in Fig. 11 where one sees the red and blue wings bracketing the central absorption in the southern part of MM1.

We searched for traces of systematic motion in the high-velocity gas seen in  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  emission. Figure 11 shows the spatial distribution of the line emission, with redshifted gas in red, blueshifted gas in blue, and intermediate velocity gas in green. The quiescent gas is seen to trace a flattened envelope which is also seen in the continuum (Fig. 4), while the high-velocity gas is seen to trace two lobes which fall on a line which is approximately perpendicular to this envelope. If these lobes are due to a bipolar outflow, Fig. 11 indicates a size of  $\sim 36''$  end-to-end or  $\sim 18''$  peak-to-peak, which at 5.5 kpc (Zhang et al. 2014) corresponds to 0.48–0.96 pc. While the map data in Fig. 11 indicate a velocity range of only  $\sim 31 \text{ km s}^{-1}$  FWZI or  $\sim 25 \text{ km s}^{-1}$  FWHM for the outflow, the pointed spectrum in Fig. 10 shows that  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  emission is detected at higher velocities, up to  $\sim 50 \text{ km s}^{-1}$  FWZI. Furthermore, Fig. 11 indicates a low degree of collimation for the outflowing gas, as often seen in single-dish observations of outflows from massive star-forming regions (Beuther & Shepherd 2005).

Outflow activity from W43-MM1 has been observed before by Sridharan et al. (2014), who used the SMA to image the CO 3–2 line emission from this region at  $5''$  resolution. Their data reveal gas shifted to velocities up to  $\sim 70 \text{ km s}^{-1}$ , which is spatially resolved into three distinct bipolar outflows, which together cover an area of  $20\text{--}30''$  in size. Thus, our observed high-velocity  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  emission may trace dense gas associated with the high-velocity CO seen with the SMA, and the apparent low collimation may be due to confusion between multiple bipolar outflows.

#### 5. Envelope absorption and velocity gradient

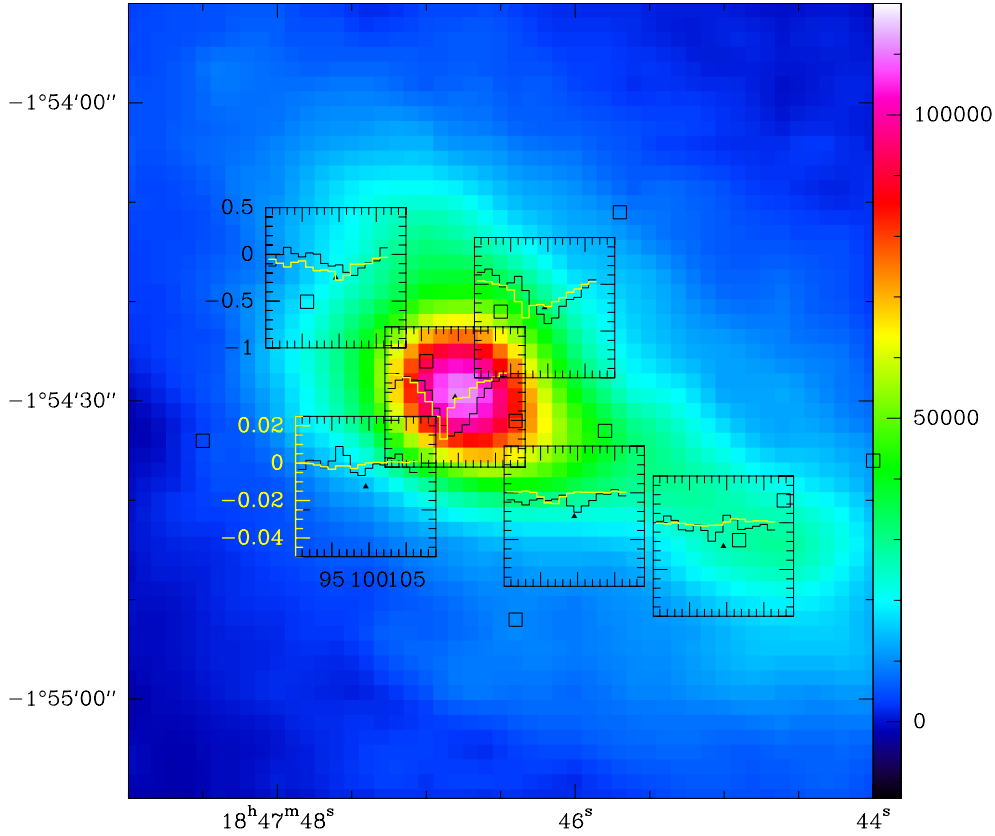
Figure 7 points to the same effect: the blue wing is seen in absorption to the SW and the red wing to the NE. The total velocity shift is around 4 to  $5 \text{ km s}^{-1}$ .

A major question is how the massive protostellar core can be observed directly. It appears that the main isotope  $p\text{-H}_2\text{O } 1_{11}-0_{00}$  absorption at 1113 GHz comes from the envelope rather than the region above 100 K (Herpin et al. 2012). In order to assess whether the  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption is due to the cool envelope or the protostellar core, we used the line-to-continuum ratio of the  $p\text{-H}_2\text{O } 1_{11}-0_{00}$  line to estimate the optical depth at 1113 GHz, implicitly assuming that all of the absorption is due to the envelope. With a minor calibration error, it is easy to underestimate or overestimate the  $p\text{-H}_2\text{O } 1_{11}-0_{00}$  optical depth. If the absorption were underestimated, one would expect to have extreme variations in the optical depth as, for example, a 99.99% absorption will yield a much higher optical depth than 99% absorption and the variations due to noise are at this level. If the absorption is overestimated, then the  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption will also be overestimated. The absorption is shown in Fig. 6, where the highest contour is at an absorption level of 78% over a width of  $4 \text{ km s}^{-1}$ , consistent with the somewhat higher absorptions we obtain for the  $2 \text{ km s}^{-1}$  width used in these calculations. Then, assuming an abundance ratio  $\text{H}_2\text{O}/\text{H}_2^{18}\text{O}$  of 322, corresponding to the galactocentric radius of W43 (Wilson & Rood 1994), we calculate the optical depth and (using the continuum flux) the expected absorption line intensity. This can be directly compared to the observed  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption spectra.

Figure 12 shows the observed  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption and predicted  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  envelope absorption spectra superposed on a SCUBA  $450 \mu\text{m}$  map. It can be seen that while there is  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption associated with the envelope, the expected line intensity is below 0.04 K everywhere. The observed  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption reaches 1 K in the central regions, incompatible with the hypothesis that the  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption is due to the envelope. In both the predicted envelope spectra and observed spectra,  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  lines are predicted/detected in regions along the major axis of W43 MM1.

Therefore, the  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  absorption is due to the central core region, which has  $T_{\text{mb}} > 100 \text{ K}$ . As also shown in Fig. 7, the velocities of the  $p\text{-H}_2\text{O } 1_{11}-0_{00}$  and  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  lines are not the same, further illustrating that they trace different components. The  $^{13}\text{CO } (10-9)$  emission (Fig. 1), tracing warm dense material and observed at the same positions with exactly the same beam size, is systematically at a less redshifted velocity than the  $p\text{-H}_2^{18}\text{O } 1_{11}-0_{00}$  implying that the water over 100 K is infalling.

Figures 5 and 6 visualize the absorbed fraction, for various velocity ranges, of the  $p\text{-H}_2\text{O } 1_{11}-0_{00}$  line. The situation in Fig. 6 is clearly different from that in Fig. 5. The diffuse line of sight cold cloud mapped in Fig. 5 does not show the same localized



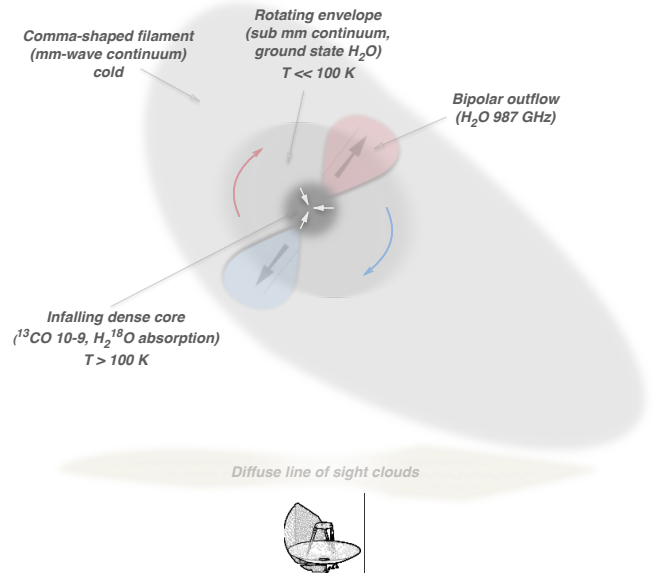
**Fig. 12.** Observed p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> spectra (black line and temperature scale) and predicted p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> (yellow line and temperature scale) due exclusively to the cool outer material believed to be responsible for the H<sub>2</sub><sup>16</sup>O absorption. The zero levels are at the same position, but there is a factor of 4 difference in the scales, showing clearly that the absorption from the cool material is far from sufficient to explain the observed absorption. Units are km s<sup>−1</sup> for the velocity and Kelvin for the line temperature. The background is a SCUBA 450 μm image with a linear scale and a peak at 110 Jy/beam. The triangles indicate the positions of the spectra. The open squares indicate the Louvet et al. (2014) sources.

structure seen in Fig. 6. The three upper panels of Fig. 6 show respectively the blue wing, the central part where the absorption is generally higher, and the red wing. There is little absorption to the NE in the blue wing and to the SW in the red, showing a velocity gradient in the envelope.

Emission/absorption of p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> is also present outside the central beam. We believe that this is due to the other massive dense cores found by Louvet et al. in W43 MM1. Their positions are indicated in Fig. 12 and this can be compared with the p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> optical depth as seen in Fig. 7, showing the same result in a slightly different manner.

## 6. Source structure and infall rate

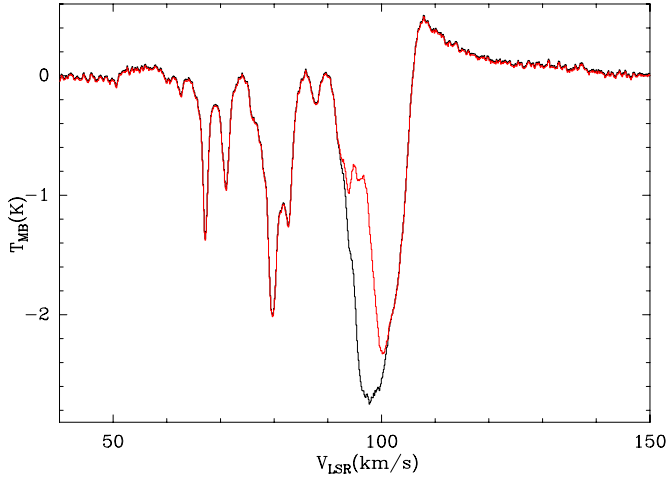
Based on results of the previous sections, Fig. 13 sketches out the kinematics and physical structure of W43 MM1; the sizes and geometry of the various components are approximations. The extended comma-shaped continuum component contains the infalling W43 MM1 core embedded in a colder envelope. The saturated p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> water line comes from the colder envelope whereas p-H<sub>2</sub><sup>18</sup>O 1<sub>11</sub>–0<sub>00</sub> and <sup>13</sup>CO (10–9) come from the hotter infalling core. The envelope component is resolved by the *Herschel* beam at 1113 GHz. Its possible rotation is suggested by the water lines velocity gradient seen around the core along the comma-shaped axis. This sketch does not show the multi-cored substructure seen in Louvet et al. (2014), and gives only the global orientation of the outflow seen in p-H<sub>2</sub>O 2<sub>02</sub>–1<sub>11</sub>.



**Fig. 13.** Various W43 MM1 components derived from the spectral maps. The arrows indicate the corresponding velocity fields (infall, outflow, and rotation). The observing direction is indicated by the bottom telescope image whose dimension gives an idea of the beam size.

### 6.1. Cold foreground material

As explained in Sect. 4 and Fig. 7, a cold foreground cloud embedding MM1 contributes to the blue side of the line profile of



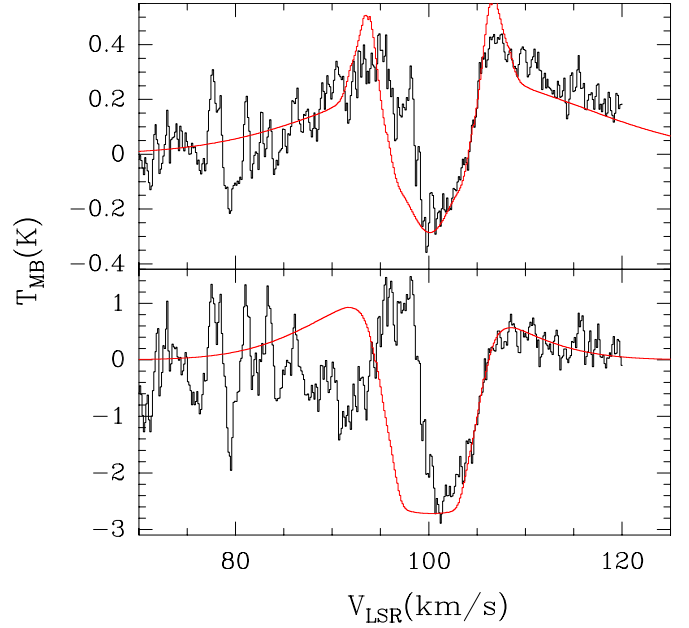
**Fig. 14.** Original HIFI spectra of the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> line in black (from Herpin et al. 2012). In red the spectra after removing of the cold component.

the para ground state line of water. We modeled this blending component using our map data as follows. Figure 5 (see left box) shows that the opacity of this material can, at first, be considered rather constant between MM1 and the southwest part of the map allowing us to make the raw hypothesis that the physical conditions of this material are similar enough both at MM1 and at a location around the offset (–21″, –14″). We fit the line at this off-core position, and we then scale this result using the continuum ratio at MM1 and at the off-core position in order to remove it from the spectrum observed towards MM1. Subtracting this absorption from the pointed spectra of MM1 from Herpin et al. (2012), we get an estimate of the *true* water line profile toward the massive dense core MM1 (see Fig. 14). The new line profile is narrower and the absorption less deep, now peaking at 100 km s<sup>–1</sup>, i.e., redshifted relative to the source velocity.

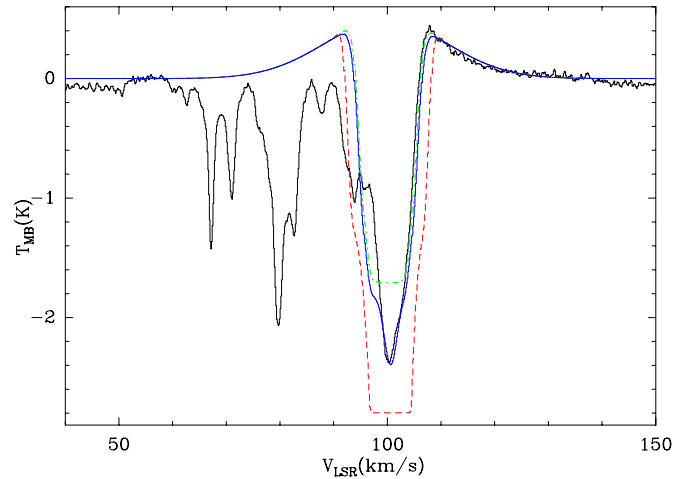
This blue side cold component is also seen in the ortho ground state lines o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and o-H<sub>2</sub>O 2<sub>12</sub>–1<sub>01</sub>. Unfortunately, missing the corresponding H<sub>2</sub>O line maps, we cannot apply the method described above for the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> line profile. Still, we can get an estimate of this blue side blending by comparing the many diffuse line of sight cold cloud components seen in o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and o-H<sub>2</sub>O 2<sub>12</sub>–1<sub>01</sub> between 0 and 90 km s<sup>–1</sup> to their p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> counterparts. We note that their integrated intensity (for both o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and o-H<sub>2</sub>O 2<sub>12</sub>–1<sub>01</sub>) can be scaled with almost the same value at all velocities from each p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> counterpart. From our hypothesis that the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> blue side material has a diffuse line of sight cold cloud nature, we can then model it for the o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and o-H<sub>2</sub>O 2<sub>12</sub>–1<sub>01</sub> main water line by scaling our fit for p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> with the scale ratio found for the o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and o-H<sub>2</sub>O 2<sub>12</sub>–1<sub>01</sub> “cold” components. We can now derive the corrected o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and o-H<sub>2</sub>O 2<sub>12</sub>–1<sub>01</sub> water line profiles (see Fig. 15).

## 6.2. Warm protostellar envelope

Considering that these new line profiles give us a more realistic view of the MM1 source, we decide to revisit the analysis of Herpin et al. (2012) who needed a cold water cloud to properly model the various line profiles. We use the same physical structure for the W43 MM1 source and first apply the same model (same foreground clouds, turbulent velocity, infall, and water



**Fig. 15.** HIFI spectra after removing of the cold component for the 557 (top) and 1669 GHz (bottom) lines. The new model is shown in red.



**Fig. 16.** New HIFI spectra of the p-H<sub>2</sub>O 1<sub>11</sub>–0<sub>00</sub> line in black (from Herpin et al. 2012). The outputs from the model with the Herpin et al. (2012) parameters, and new parameters without and with a foreground cloud are shown respectively in red, green, and blue.

abundances) as Herpin et al. (2012). Figure 16 shows that the line profile is not well reproduced: the modeled spectra are too broad and exhibit a saturated deep absorption. The best model is obtained by only decreasing the global outer water abundance from  $8 \times 10^{-8}$  to  $8 \times 10^{-9}$ ; the inner water abundance and the turbulence velocity increasing with radius have not changed. Surprisingly, a weaker cold component (with an opacity of 0.3) is still needed to reach the absorption dip. The same conclusions apply to the o-H<sub>2</sub>O 1<sub>10</sub>–1<sub>01</sub> and o-H<sub>2</sub>O 2<sub>12</sub>–1<sub>01</sub> lines (opacity of 0.4 for the cold component, see Fig. 15).

The outer water abundance in other mid-IR quiet high-mass protostellar objects (HMPO) has been found by Herpin et al. (2016) and Marseille et al. (2010) to be a few  $10^{-8}$  (but  $1.4 \times 10^{-7}$  for DR21(OH)); however, other estimates in some more evolved massive objects lead to lower values:  $2.4$ – $6.3 \times 10^{-9}$  in AFGL2591 (Choi et al. 2015) and  $1$ – $5 \times 10^{-9}$  for some bright HMPOs, hot molecular cores, and Ultra Compact HII regions



(Choi 2015). In addition, low-mass protostellar objects tend to have an outer abundance of  $\sim 10^{-8}$  (Van Dishoeck et al. 2011) or even lower ( $3 \times 10^{-10}$ , Mottram et al. 2013). Moreover, Snell et al. (2000) or Caselli et al. (2012) have estimated very low abundances ( $10^{-10}$ – $10^{-8}$ ) in cold regions. Hence the value derived in our study is compatible with that found in cold outer regions. In the cold envelopes, water is mainly found as ice and can sublime through far-ultraviolet (FUV) photodesorption (photons from the interstellar radiation field, ISRF, and from cosmic rays interacting with  $H_2$ , see Schmalzl et al. 2014). Of course, part of this water vapour can then be photodissociated through FUV photons. Variations of the ISRF (e.g., proximity of a nearby bright star) or self-shielding, for instance, can lead to different outer abundances for different objects.

While we find lower outer water abundances than derived from pointed observations, it is unlikely that this can be generalized to other sources. W43 MM1 is indeed part of a complex cluster embedded in an elongated structure (Louvet et al. 2014). Only further studies of similar maps observed toward these objects can help to confirm or revise previous water abundances.

In the new models, as for  $NH_3\ 3_{2+}-2_{2-}$  in Wyrowski et al. (2012), the lines are now centered at  $100\text{ km s}^{-1}$  rather than at  $98.6\text{ km s}^{-1}$ , as used in Herpin et al. (2012), hence redshifted relative to the source velocity, which is indicative of infall. While in Herpin et al. (2012), the ground-state lines were dominated by the cold cloud (in absorption), and thus did not show any infall signature, these new water line profiles do indicate infall. This infall signature is now well reproduced by this new model using the same infall velocity that Herpin et al. (2012) used ( $-2.9\text{ km s}^{-1}$ ), then mainly derived from the 987 and 752 GHz water lines. Before subtraction of the widespread cool component, the low-energy  $H_2O$  lines were not sensitive to infall, which could only be traced at higher energies and/or with isotopologues. Now the  $H_2O$  lines yield the same level of infall. The main consequence of this new outer water abundance is that the water content in the outer region is an order of magnitude lower than previously estimated:  $3.7 \times 10^{-4} M_\odot$  (instead of  $3 \times 10^{-3} M_\odot$ ).

## 7. Conclusions

The main conclusion from this work is that map data are an essential complement to pointed observations to ensure reliable estimation of physical and chemical parameters from submillimeter line spectra. In particular, molecular abundances derived only from pointed spectra may be off by a factor of  $\approx 10$ , while estimates of turbulent velocities and infall/outflow rates are less affected.

Our second conclusion is that lines of  $H_2^{18}O$ , as well as excited-state  $H_2O$  lines, trace the inner envelopes of high-mass protostars where abundances are enhanced due to grain mantle evaporation and/or high-temperature gas-phase chemistry. However, deriving these abundances requires careful subtraction of the contribution from the outer envelope.

Third, we find that warm chemically enriched gas in W43 is not confined to the MM1 position, but is also found towards other cores seen in warm dust emission by Louvet et al. (2014).

Finally, the HIFI maps of the W43 region reveal velocity gradients in the envelope of the MM1 core as well as in the surrounding gas, which are indicative of rotation. Figure 13 summarizes the global picture derived for the MM1 region.

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